



Equivalent Multiple Complex SUSY For Real SUSY

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Author's contribution

The sole author designed, analyzed and interpreted and prepared the manuscript.

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ABSTRACT

We notice real *SUSY* Hamiltonians have multiple equivalent complex Hamiltonians which may be (i) \mathcal{PT} invariant (ii) \mathcal{T} invariant or (iii) combination of both in nature . These three types of complex Hamiltonians give the same energy spectrum . We present here analytical results for the exactly solvable system and numerical results for others.

Keywords: Supersymmetry; \mathcal{PT} symmetry; real spectra; complex hamiltonians.

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1 INTRODUCTION

Our understanding on real spectra in quantum physics has been drastically changed after the thought breaking idea of Bender and Boettcher [1], who introduced the concept of \mathcal{PT} symmetry . The operator \mathcal{P} stands for parity ,reflecting the behaviour $x \rightarrow -x ; p \rightarrow -p$ and $i \rightarrow i$. Similarly the operator \mathcal{T} represents time reversal ,reflecting the behaviour $x \rightarrow x ; p \rightarrow -p$ and $i \rightarrow -i$. However , \mathcal{PT} symmetry understanding took a new turn when Jones and Mateo [2],theoretically proved that inverted quartic oscillator

$$H = p^2 - x^4 \quad (1)$$

has equivalent hermitian operator reflecting the iso-spectral character . Later on Nanayakkara and Mathanaranjan [3]noticed that one dimensional complex Hamiltonian

$$H = p^2 - x^4 + 4ix \quad (2)$$

also posseses equivalent hermitian operator reflecting iso-spectra . One simple question comes to mind that whether real Hamiltonians have complex counter part reflecting iso-spectral behaviour ? In order to address this question we consider supersymmetry as an ideal example .

2 REAL SUSY

Here, we simply consider, SUSY Hamiltonians in short as follows . The generated Hamiltonians in terms of superpotential W , can be written as [4-7]

$$H^+ = p^2 + \frac{dW(x)}{dx} + W^2 \quad (3)$$

and

$$H^- = p^2 - \frac{dW(x)}{dx} + W^2 \quad (4)$$

For SUSY energy conditions

$$E_n^{(+)} = E_{n+1}^{(-)} \quad (5)$$

with

$$E_0^{(-)} = 0 \quad (6)$$

let us consider two quadratic exactly solvable Hamiltonians as

$$H^- = p^2 + x^2 - 1 \quad (7)$$

$$H^+ = p^2 + x^2 + 1 \quad (8)$$

We can have another two Hamiltonians as

$$H^- = p^2 + x^6 + 2x^4 - 2x^2 - 1 \quad (9)$$

$$H^+ = p^2 + x^6 + 2x^4 + 4x^2 + 1 \quad (10)$$

The computed eigenvalues are tabulated in table 1 using matrix diagonalisation method [8]. Further ,for the iso-spectral energy condition

$$E_n^{(+)} = E_n^{(-)} \quad (11)$$

we can have the Hamiltonians as,

$$H^- = p^2 + x^4 + x^2 - 2x + 0.25 \quad (12)$$

$$H^+ = p^2 + x^4 + x^2 + 2x + 0.25 \quad (13)$$

3 COMPLEX SUSY

Before, going to introduce complex SUSY , we would like to bring the attention of reader, an interesting idea on complex transformation of momentum [9-12] in terms of co-ordinate as

$$p \rightarrow p + ix \quad (14)$$

It has been explicitly addressed in the case of Harmonic Oscillator[9-12] . In this paper we apply the same to real SUSY operators as

$$H_1^+ = p^2 - x^2 + i(xp + px) + \frac{dW(x)}{dx} + W^2 \quad (15)$$

and

$$H_1^- = p^2 - x^2 + i(xp + px) - \frac{dW(x)}{dx} + W^2 \quad (16)$$

The above two Hamiltonians are \mathcal{PT} invariant in nature . In our view ,the above two Hamiltonians must retain iso-spectral behaviour with that of real SUSY Hamiltonians .Now consider another transformation as

$$p \rightarrow p + i \quad (17)$$

The new complex Hamiltonians are as follows

$$H_2^+ = p^2 + 2ip - 1 + \frac{dW(x)}{dx} + W^2 \quad (18)$$

and

$$H_2^- = p^2 + 2ip - 1 - \frac{dW(x)}{dx} + W^2 \quad (19)$$

The above two Hamiltonians are \mathcal{T} invariant in nature . Now consider combination of these two and write two new Hamiltonians as

$$H_3^+ = p^2 - x^2 + i(xp + px) + 2ip - 2x - 1 + \frac{dW(x)}{dx} + W^2 \quad (20)$$

and

$$H_3^- = p^2 - x^2 + i(xp + px) + 2ip - 2x - 1 - \frac{dW(x)}{dx} + W^2 \quad (21)$$

Interestingly in this case the Hamiltonians are neither \mathcal{PT} invariant nor \mathcal{T} invariant in nature . In order to show explicitly we consider few cases as given below.

3.1 Complex SUSY: Analytical Result

Here we would like to state that quadratic operator can be addressed analytically.Let us discuss few lines on analytical expression for energy level relating to quadratic Hamiltonian [10,11]

$$H = h_{11}p^2 + ih_{12}(xp + px) + h_{22}x^2 + ih_1p + h_2x \quad (22)$$

having eigenvalue

$$\epsilon_n = [\sqrt{(h_{11}h_{22} + h_{12}^2)}](2n + 1) + \frac{(h_1^2h_{22} - h_2^2h_{11} - 2h_1h_2h_{12})}{4(h_{11}h_{22} + h_{12}^2)} \quad (23)$$

Here we suggest two different complex Hamiltonians as follows Now consider complex SUSY on exactly solvable real systems as:

$$H_1^+ = p^2 + i(xp + px) + 1 \quad (24)$$

$$H_1^- = p^2 + i(xp + px) - 1 \quad (25)$$

$$H_2^+ = p^2 + x^2 + 2ip \quad (26)$$

$$H_2^- = p^2 + x^2 + 2ip - 2 \quad (27)$$

$$H_3^+ = p^2 + i(xp + px) + 2ip - 2x \quad (28)$$

$$H_3^- = p^2 + i(xp + px) + 2ip - 2x - 2 \quad (29)$$

Using the above expression ,one can see that

$$H_3^-, H_2^-, H_1^- = 2n \quad (30)$$

and

$$H_3^+, H_2^+, H_1^+ = 2n + 2 \quad (31)$$

Here , n=0,1,2,3 Interested reader can easily verify the SUSY energy conditions.

3.2 Complex SUSY: Numerical Result

Here we consider the complex SUSY Hamiltonians as

$$H_1^- = p^2 + x^6 + 2x^4 - 3x^2 - 1 + i(xp + px) \quad (32)$$

$$H_1^+ = p^2 + x^6 + 2x^4 + 3x^2 + 1 + i(xp + px) \quad (33)$$

$$H_2^- = p^2 + x^6 + 2x^4 - 2x^2 + 2ip - 2 \quad (34)$$

$$H_2^+ = p^2 + x^6 + 2x^4 + 4x^2 + 2ip \quad (35)$$

$$H_3^- = p^2 + x^6 + 2x^4 - 3x^2 - 2 + i(xp + px) + 2ip - 2x \quad (36)$$

$$H_3^+ = p^2 + x^6 + 2x^4 + 3x^2 + i(xp + px) + 2ip - 2x \quad (37)$$

The above Hamiltonians can not be solved analytically. For numerical results we apply matrix diagonalisation method [8] as follows

$$H|\Psi\rangle = E|\Psi\rangle \quad (38)$$

where

$$|\Psi\rangle = \sum_m A_m |m\rangle \quad (39)$$

In the above $|m\rangle$ is the harmonic oscillator wave function which satisfies the eigenvalue relation

$$(p^2 + x^2)|m\rangle = (2m + 1)|m\rangle \quad (40)$$

Further in general ,for SUSY Hamiltonian we get nine term recurrence relation as

$$A_{m-6}P_m + A_{m-4}Q_m + A_{m-2}R_m + A_{m-1}S_m + A_m T_m + A_{m+1}U_m + A_{m+2}V_m + A_{m+4}W_m + A_{m+6}Y_m = 0 \quad (41)$$

Where

$$P_m = \langle m|H|m-6\rangle \quad (42)$$

$$Q_m = \langle m|H|m-4\rangle \quad (43)$$

$$R_m = \langle m|H|m-2\rangle \quad (44)$$

$$S_m = \langle m|H|m-1\rangle \quad (45)$$

$$U_m = \langle m|H|m+1\rangle \quad (46)$$

$$V_m = \langle m|H|m+2\rangle \quad (47)$$

$$W_m = \langle m|H|m+4\rangle \quad (48)$$

$$Y_m = \langle m|H|m+6\rangle \quad (49)$$

$$T_m = \langle m|H|m\rangle - E \quad (50)$$

For the benefit of readers we present diagonal elements as given below .

$$\langle m|H^+|m\rangle = 2.5m^3 + 6.75m^2 + 13m + 6.875 \quad (51)$$

$$\langle m|H_1^+|m\rangle = 2.5m^3 + 6.75m^2 + 12m + 6.375 \quad (52)$$

$$\langle m|H_2^+|m\rangle = 2.5m^3 + 6.75m^2 + 13m + 5.875 \quad (53)$$

$$\langle m|H_3^+|m\rangle = 2.5m^3 + 6.75m^2 + 12m + 5.375 \quad (54)$$

$$\langle m|H_1^-|m\rangle = 2.5m^3 + 6.75m^2 + 6m + 1.375 \quad (55)$$

$$\langle m|H_2^-|m\rangle = 2.5m^3 + 6.75m^2 + 7m + 0.875 \quad (56)$$

$$\langle m|H_3^-|m\rangle = 2.5m^3 + 6.75m^2 + 6m + 0.375 \quad (57)$$

$$\langle m|H^-|m\rangle = 2.5m^3 + 6.75m^2 + 7m + 1.875 \quad (58)$$

In table 1 , we reflect eigenvalues along with the real SUSY Hamiltonians . Here $H_{1,2,3}^- \rightarrow E_n^{-C}$ and $H_{1,2,3}^+ \rightarrow E_n^{+C}$

Table 1. Eigenvalues of real and complex SUSY hamiltonians

n	$E_n^{(-R)}$	$E_n^{(+R)}$	E_n^{-C}	E_n^{+C}
0	0	3.373 001 0	0	3.373 001 0
1	3.373 001 0	8.743 633 3	3.373 001 0	8.743 633 3
2	8.743 633 3	15.261 907 1	8.743 633 3	15.261 907 1
3	15.261 907 1	22.749 693 9	15.261 907 1	22.749 693 9

3.3 Iso- Spectral Complex Hamiltonians: Numerical Result

Now we consider iso-spectral nature of complex SUSY Hamiltonians . The Hamiltonians considered here as

$$H_1^- = p^2 + x^4 + i(xp + px) - 2x + 0.25 \quad (59)$$

$$H_1^+ = p^2 + x^4 + i(xp + px) + 2x + 0.25 \quad (60)$$

$$H_2^- = p^2 + x^4 + x^2 + 2ip - 2x - 0.75 \quad (61)$$

$$H_2^+ = p^2 + x^4 + x^2 + 2ip + 2x - 0.75 \quad (62)$$

$$H_3^- = p^2 + x^4 + i(xp + px) - 4x + 2ip - 0.75 \quad (63)$$

$$H_3^+ = p^2 + x^4 + i(xp + px) + 2ip - 0.75 \quad (64)$$

Here we calculate energy eigenvalues using matrix diagonalisation , on solving the eigenvalue relation as stated above . Here we solve a seven term recurrence relation as given below

$$A_{m-4}Q_m + A_{m-2}R_m + A_{m-1}S_m + A_mT_m + A_{m+1}U_m + A_{m+2}V_m + A_{m+4}W_m = 0 \quad (65)$$

where

$$Q_m = \langle m|H|m - 4 \rangle \quad (66)$$

$$R_m = \langle m|H|m - 2 \rangle \quad (67)$$

$$S_m = \langle m|H|m - 1 \rangle \quad (68)$$

$$U_m = \langle m|H|m + 1 \rangle \quad (69)$$

$$V_m = \langle m|H|m + 2 \rangle \quad (70)$$

$$W_m = \langle m|H|m + 4 \rangle \quad (71)$$

$$T_m = \langle m|H|m \rangle - E \quad (72)$$

For the benefit of readers we present diagonal elements as given below .

$$\langle m|H^+|m \rangle = 1.5m^2 + 3.5m + 2 \quad (73)$$

$$\langle m|H_1^+|m \rangle = 1.5m^2 + 2.5m + 1.5 \quad (74)$$

$$\langle m|H_2^+|m \rangle = 1.5m^2 + 3.5m + 1 \quad (75)$$

$$\langle m|H_3^+|m \rangle = 1.5m^2 + 2.5m + 0.5 \quad (76)$$

$$\langle m|H_1^-|m \rangle = 1.5m^2 + 2.5m + 1.5 \quad (77)$$

$$\langle m|H_2^-|m \rangle = 1.5m^2 + 3.5m + 1 \quad (78)$$

$$\langle m|H_3^-|m \rangle = 1.5m^2 + 2.5m + 0.5 \quad (79)$$

$$\langle m|H^-|m \rangle = 1.5m^2 + 3.5m + 2 \quad (80)$$

In table 2, we reflect eigenvalues along with the real iso-spectral Hamiltonians using matrix diagonalisation method as described earlier . Here $H_{1,2,3}^- \rightarrow E_n^{-C}$ and $H_{1,2,3}^+ \rightarrow E_n^{+C}$

Table 2. Eigenvalues of real and complex Iso-Spectral hamiltonians

n	$E_n^{(-R)}$	$E_n^{(+R)}$	E_n^{-C}	E_n^{+C}
0	1.277 243 8	1.277 243 8	1.277 243 8	1.277 243 8
1	4.771 390 0	4.771 390 0	4.771 390 0	4.771 390 0
2	8.812 448 7	8.812 448 7	8.812 448 7	8.812 448 7
3	13.333 679 9	13.333 679 9	13.333 679 9	13.333 679 9

4 CONCLUSION

In this paper we have found that for real *SUSY* Hamiltonians, there are multiple complex equivalent Hamiltonians reflecting iso-spectral behaviour. Hence we believe all bounded operators are associated with equivalent complex operators i.e

$$\begin{array}{c} \text{Bounded Operator} \\ \updownarrow \\ \overbrace{\mathcal{PT} + \mathcal{T} + \mathcal{PT} \sim \mathcal{T}} \end{array}$$

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COMPETING INTERESTS

Author has declared that no competing interests exist.

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