

Journal of Materials Science Research and Reviews

Volume 7, Issue 2, Page 204-209, 2024; Article no.JMSRR.116184

Optimizing Power Factor in SixGe1-x Alloys and Some Other Thermoelectric Materials: Insights and Empirical Relationships

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Authors' contributions

This work was carried out in collaboration between both authors. Both authors read and approved the final manuscript.

Article Information

Open Peer Review History:

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Short Communication

Received: 23/02/2024 Accepted: 27/04/2024 Published: 01/05/2024

ABSTRACT

In this article, the issue of determining the power factor maximums of Si_xGe_{1-x} alloy with different compositions is investigated. This issue is discussed also for other thermoelectrics based on literature data. It is shown that for the values of the Seebeck coefficient in the interval (1-4)10⁻⁴V/K, (PF)_{max} corresponds to the minimum of the specific electrical conductivity. The interdependence of these thermoelectric parameters has a regular character. The dependences $Ig(\sigma S^2)_{max} - Ig\sigma_{min}$ for various thermoelectrics based on literature and our data in the corresponding interval of changes of variables are described by a single empirical expression $Ig(\sigma S^2)_{max} \cong 0.583(Ig\sigma_{min})^2 \cdot 3.332Ig\sigma_{min}$. The temperature dependence of the Seebeck coefficient described by the equation $S=\frac{3}{2}\frac{k_B}{q_e}InT+C$, where C depends on the concentration and effective bmass of charge carriers, as well as on the Debye temperature. This dependence has the same character as dependence $\sigma S^2 - T$ and have

maximum around 1100K.

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Additionally, the dependences of power factor on specific electrical conductivity were studied. For different thermoelectrics they are described by empirical expressions such as $\sigma S^2 \cong (a/\sigma)$ -b or $\sigma S^2 \cong c\sigma^{-d}$, where constants a, b, c and d>0. For a number of thermoelectrics, the dependences σS^2 on the Seebeck coefficient are rectilinear, while for another series they are power-law. And this is a more general case. The temperature dependences of σS^2 and of scaled power factor (Bs) are also studied. Both are described by a quadratic equation. It has been established that the maximum power factor corresponds to the minimum Bs. Thus, in this case to estimate (σS^2)_{max} can be used as σ_{min} , and (Bs)_{min}.

Keywords: Thermoelectrics; power factor; electrical conductivity.

1. INTRODUCTION

Since the seventies of the last century, the Si_xGe_{1-x} alloy has been used in thermoelectric generators in spacecraft for long-distance flights with a long service life [1-3]. This alloy is also used in many branches of science and technology [4-15].

N-type Si_xGe_{1-x} has a number of advantages over the P-type: the maximum of figure of merit (ZT) is about 2.5 times greater for N-type than for P-type (at concentration of charge carriers n= 3.2.10²⁶m⁻³ and at the same compositions). This follows from the fact that the specific electrical conductivity is 2.5-3 times higher, Seebeck coefficient is 1.4-2 times larger (accordingly, the power factor is ~2 orders of magnitude greater), and thermal conductivity coefficient 1.1-1.3 times is smaller at the same temperatures. Also with ≥10¹⁹ and fluencies neutron irradiation temperatures ≥600°C N-type Si_xGe_{1-x} is more radiation resistant [16].

This article is a continuation of work [17], where the maximization of the figure of merit of thermoelectric material SiGe is considered. Together with ZT, power factor (PF) is an important energy characteristic of thermoelectrics, which is included in the expression of the ZT: $PF \equiv \sigma S^2$, where σ is the specific electrical conductivity, S - Seebeck coefficient. Here we consider N-SixGe1-x alloy (x=0.7, 0.72, 0.76, 0.8 and 0.83) based on experimental data of [18]. Since the remaining thermoelectric materials are also taken from the literature, the experimental section is not included.

2. DISCUSSION

According to a fairly large number of works, the range of changes of the Seebeck coefficient is $(1-4)10^{-4}$ V/K. With such a relatively narrow

range, the σS^2 - S dependence is almost rectilinear: $\sigma S^2 = kS + b$, where k is the slope of the straight lines, and b is the ordinate of the point of their intersection with the axis σS^2 when extrapolating these lines to S \rightarrow 0. From the last equation we get: $S = \frac{k}{2\sigma} + \left[\left(\frac{k}{2\sigma}\right)^2 + \frac{b}{\sigma}\right]^{1/2}$. After simple transformations we will have: $\sigma S^2 = \frac{k}{2\sigma} + b$, that is, in this case (PF)_{max} will correspond to σ_{min} .

Fig 1 shows σS^2 - S dependencies for different thermoelectric materials, namely for a, b and c. These dependencies are constructed based on experimental data from the following works: [18] (SiGe), [19] (Bi₂Sr_{1.925}[Sr(BO₂)₂]_{0.075}CO_{1.8}O_y, Bi₂Sr₂CO_{1.8}O_y), [20] (SiGeMo_{0.2}), and [21] (Sb_{1.7}Bi_{0.4}Te₃).

Fig 2 shows $lg(\sigma S^2)_{max} - lg\sigma_{min}$ dependences for different thermoelectrics using literature and our data. In the obtained intervals of change of variables, these dependencies can be described as a whole by the empirical expression $lg(\sigma S^2)_{max} \cong 0.583(lg\sigma_{min})^2 - 3.332lg\sigma_{min}$ (solid line in Fig.2). The entire dependence is divided into three areas: (I) $10^{-2} < (PF)_{max} < 10^{-3}$, (II) $10^{-3} <$ $(PF)_{max} < 10^{-4}$ and (III) $10^{-4} < (PF)_{max} < 10^{-5}$ W/K²m. By means of this approximate relationship, it will be possible to predict the order of $(PF)_{max}$ according to σ_{min} . The scatter of points around the averaged line is apparently due to the non-strict rectilinearity of the corresponding $\sigma~S^2$ - S dependences. When these dependencies are practicalli rectilinear, then the points fit well on the curve (compare Figs.1 and 2).

Fig 3 shows the temperature dependences of $\sigma S^2 - T$ and S - InT for Si_{0.72}Ge_{0.28} and Si_{0.76}Ge_{0.24} (dependencies for other compositions have the same form). From this figure it is clear that these dependences have a maximum at the same temperature (~1100K).

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Fig. 1. Typical dependences σS^2 - S: (a) Si_xGe_{1-x}, x=0.7 (o), 0.72 (Δ), 0.76 (\Box), 0.8 (\diamond) and 0.83 (∇); (b) (o) Bi₂Sr_{1.925}[Sr(BO₂)₂]_{0.075}Co_{1.8}O_y, (Δ) Bi₂Sr₂Co_{1.8}O_y, \Box - SiGeMo_{0.2}, and (\diamond) Sb_{1.7}Bi_{0.4}Te₃. [σS^2]=W/K²m, [S]=V/K



Fig. 2. Dependences Ig(σS²)_{max} – Igσ_{min}. The coordinates of the points are calculated according to the data of following works: sect.I - [18] (Si_{0.7}Ge_{0.3}), [20] (SiGe, SiGeMo_{0.2}), PbTe [22]; sect.II - [21] (Sb_{1.7}Bi_{0.4}Te₃), [23] Bi₂Te₃, (Bi_{0.98}Sn_{0.02})₂Te_{2.7}Se_{0.3}), [24] Si_{0.68}Ge_{0.32}Ga, [25] Ni₂CuCrFe); sect.III - [19] Bi₂Sr₂Co_{1.8}O_y, Bi₂Sr_{1.9}[Sr(BO₂)₂]_{0.1}Co_{1.8}O_y (and the same thermoelectrics with different composition of components). Figurative points correspond to Fig.1: Δ - Bi₂Sr₂Co_{1.8}O_y, □ - Si_{0.72}Ge_{0.28}, ◊ - Bi₂Te₃, ∇, ○ - SiGeGa, △ - SiGeMo_{0.2}

It should be noted that for practically all other thermoelectrics considered above (and also for $P-Si_xGe_{1-x}$), a rectilinear dependence of the Seebeck coefficient on the natural logarithm of

temperature was observed (Fig.3). Therefore, for them the following formula was used [25,26]:

$$S = \frac{3}{2} \frac{k_B}{q_e} \ln T + C, \qquad (1)$$

where C depends on the concentration and effective mass of charge carriers, as well as on the Debye temperature (k_B - Boltzmann's constant, qe - elementary charge, T - absolute temperature)^(*). Fig 3(b) shows these dependences for Sb_{1.7}Bi_{0.4}Te₃, Tl₉Sb_{0.99}Sn_{0.1}Te₆ [27] and SiGeMo_{0.2}. From Fig.4(a) it is clear that the experimental points form a practically single set with an average overall slope $K \equiv tg\alpha \cong 1.122 \cdot 10^{-4}$. The slopes of the lines presented in Fig.3(b) are tgα ≅ 1.324, 1.475 and $1.050 \cdot 10^{-4}$, respectively. These slopes which approach to $\frac{3 k_B}{2} \cong 1.293 \cdot 10^{-4}$ with more K= or less 2 qe accuracy.



Fig. 3. Dependencies $\sigma S^2 - T$ and S - InT for $Si_{0.72}Ge_{0.28}$ (o, Δ) and $Si_{0.76}Ge_{0.24}(\bullet, \blacktriangle)$. (In1100 \cong 7, the scales are different on the abscissa axes) [σS^2]=W/K²m, [S]=V/K, [T]=K



Fig. 4. Dependences S – InT for Sb_{1.7}Bi_{0.4}Te₃ (o), TI₉Sb_{0.99}Sn_{0.1}Te₆ (Δ) and SiGeMo_{0.2} (\Box). A straight line without points was constructed at tg $\alpha \cong 1.293 \cdot 10^{-4}$ with arbitrarily taken b=-4.4. [S]=V/K, [T]=K

Since the values of K do not coincide exactly $\frac{3}{2}\frac{k_B}{a}$, it is difficult to estimate the with 2 qe corresponding values of the constant A (see Footnote below) from the values of C calculated below. However. according to as vet unpublished data for P-Sio 7Geo 3. the temperature dependence of mobility is described by the expression $\mu\cong\frac{1}{2}\,T^{\text{-3/2}},$ which indicates phonon scattering of charge carriers [28].

NOTE

 ${}^{(*)}C = \frac{k_B}{q_e} \Big[A + \ln \frac{2(2\pi m^* k_B)^{3/2}}{(2\pi h^3)^3 n} \Big], \text{ where A depends on the scattering mechanism and takes values from 2 to 4 (m^* and n - effective mass and concentration of charge carriers). }$

3. CONCLUSION

A method is proposed for determining the maximum power factor of Si_xGe1_{-x} alloys of different compositions and different conductivities. Other thermoelectrics are also considered based on literature data.

The dependences $\lg(\sigma S^2)_{max} - \lg\sigma_{min}$ for various thermoelectrics based on literature and our data in the corresponding interval of changes of variables are described by a one empirical expression $\lg(\sigma S^2)_{max} \cong 0.583(\lg\sigma_{min})^2$ -3.332 $\lg\sigma_{min}$.

The temperature dependence of the Seebeck coefficient described by the equation $S = \frac{3}{2} \frac{k_B}{q_e} \ln T + C$. This dependence has the same character as dependence $\sigma S^2 - T$ and have maximums around 1100K.

COMPETING INTERESTS

Authors have declared that no competing interests exist.

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